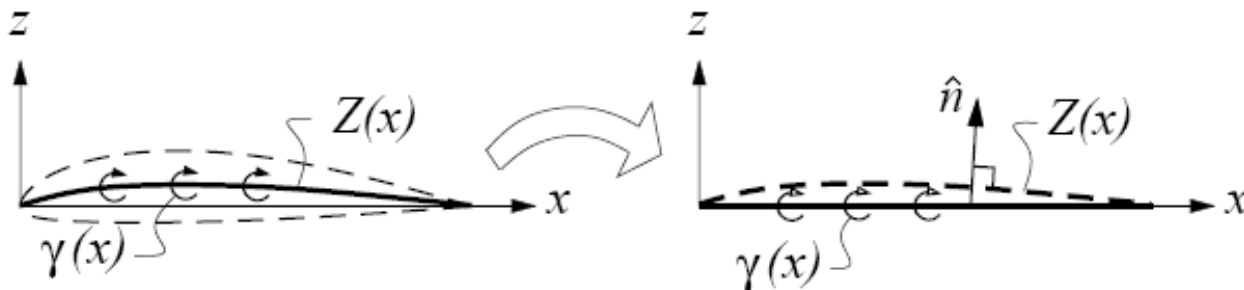


## < 4.7 Classical Thin Airfoil Theory >

### ❖ The Symmetric Airfoil

#### \* Assumptions

- i) The camber line is one of the streamlines
- ii) Small maximum camber and thickness relative to the chord
- iii) Small angle of attack



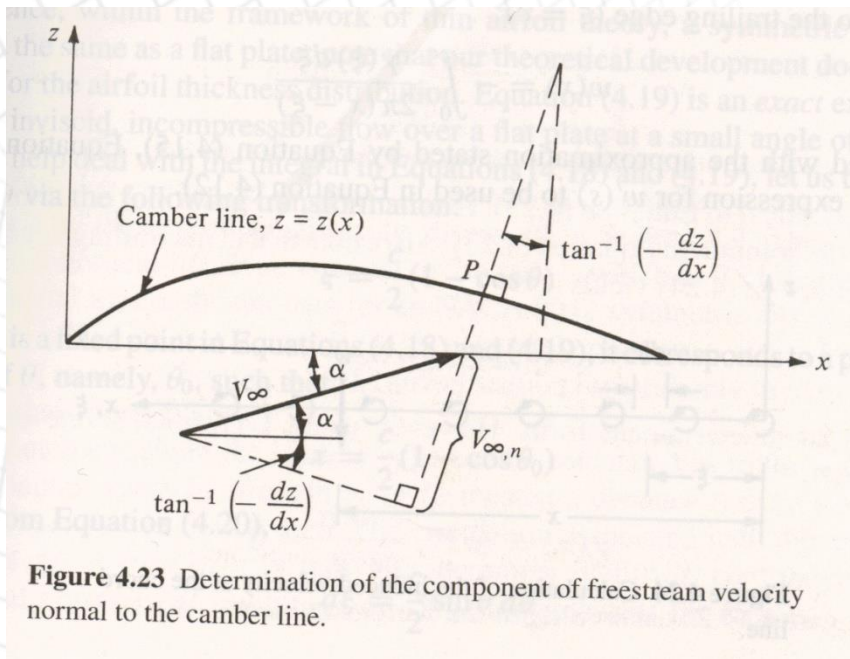
#### \* Purposes

- i) Find  $\gamma(s)$
- ii) Use Kutta-Joukowski theorem,  $L' = \rho V \Gamma$

# Incompressible Flow over Airfoils

## < 4.7 Classical Thin Airfoil Theory >

### ❖ The Symmetric Airfoil



\* The component of free-stream velocity normal to the mean camber line at P

$$\rightarrow V_{\infty,n} = V_\infty \sin\left(\alpha + \tan^{-1}\left(-\frac{dz}{dx}\right)\right)$$

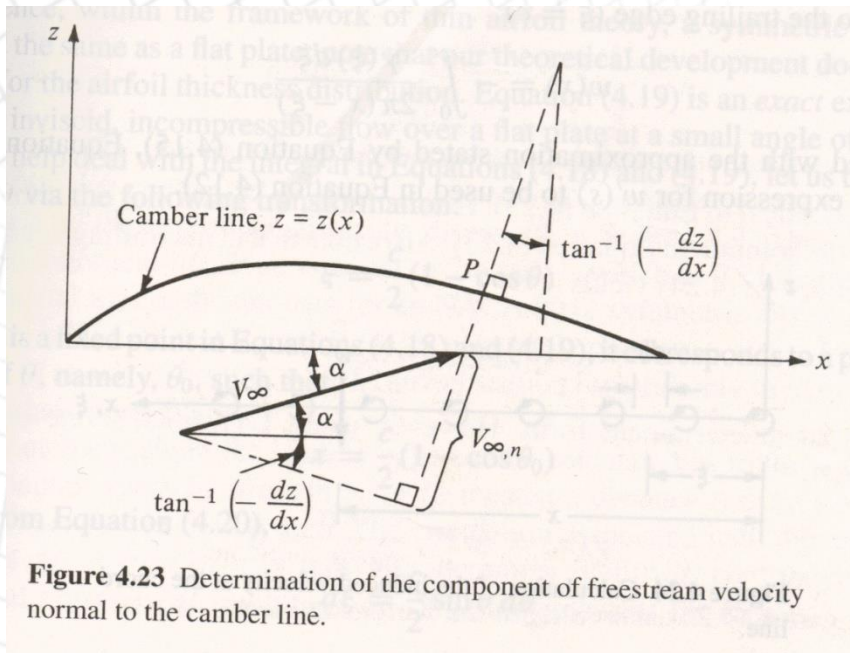
From small angle assumption

$$\rightarrow V_{\infty,n} = V_\infty \left(\alpha - \frac{dz}{dx}\right)$$

# Incompressible Flow over Airfoils

## < 4.7 Classical Thin Airfoil Theory >

### ❖ The Symmetric Airfoil



\* If the airfoil is thin,  $w'(s) \approx w(x)$

$w'(s)$  : velocity normal to the camber line  
induced by the vortex sheet

$w(x)$  : velocity normal to the chord line  
induced by the vortex sheet

\* The velocity at point  $x$  by the elemental vortex at point  $\xi$

$$\rightarrow dw = - \frac{\gamma(\xi) d\xi}{2\pi(x - \xi)}$$

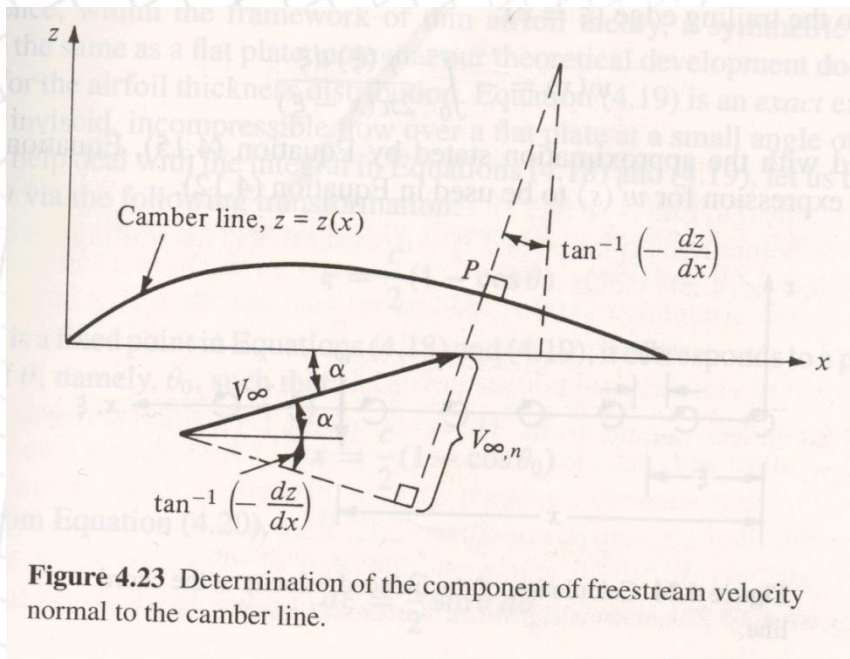
\* The velocity at point  $x$  by all the elemental vortices along the chord line

$$\rightarrow w(x) = - \int_0^c \frac{\gamma(\xi) d\xi}{2\pi(x - \xi)}$$

# Incompressible Flow over Airfoils

## < 4.7 Classical Thin Airfoil Theory >

### ❖ The Symmetric Airfoil



\* The sum of the velocity components normal to the surface at all point along the vortex sheet is zero

$$\rightarrow V_{\infty,n} + w(x) = 0$$

$$\rightarrow V_{\infty} \left( \alpha - \frac{dz}{dx} \right) - \int_0^c \frac{\gamma(\xi) d\xi}{2\pi(x-\xi)} = 0$$

$$\rightarrow \frac{1}{2\pi} \int_0^c \frac{\gamma(\xi) d\xi}{(x-\xi)} = V_{\infty} \left( \alpha - \frac{dz}{dx} \right)$$

↳ The fundamental equation of thin airfoil theory

# Incompressible Flow over Airfoils

## < 4.7 Classical Thin Airfoil Theory >

### ❖ The Symmetric Airfoil

\* Symmetric airfoil → no camber,

$$\boxed{\frac{1}{2\pi} \int_0^c \frac{\gamma(\xi) d\xi}{(x-\xi)} = V_\infty \left( \alpha - \frac{dz}{dx} \right)} \xrightarrow{\frac{dz}{dx} = 0} \boxed{\frac{1}{2\pi} \int_0^c \frac{\gamma(\xi) d\xi}{x-\xi} = V_\infty \alpha}$$

\* Transform variable  $\xi$  into  $\theta$

$$\xi = \frac{c}{2}(1 - \cos\theta) \quad \begin{cases} \theta = 0, & \xi = 0 \\ \theta = \pi, & \xi = c \end{cases}, \quad x = \frac{c}{2}(1 - \cos\theta_0), \quad d\xi = \frac{c}{2} \sin\theta d\theta$$

$$\frac{1}{2\pi} \int_0^c \frac{\gamma(\xi) d\xi}{x-\xi} = V_\infty \alpha \rightarrow \frac{1}{2\pi} \int_0^\pi \frac{\gamma(\theta) \sin\theta d\theta}{\cos\theta - \cos\theta_0} = V_\infty \alpha \rightarrow \boxed{\gamma(\theta) = 2\alpha V_\infty \left( \frac{1 + \cos\theta}{\sin\theta} \right)}$$

$$\left( \int_0^\pi \frac{n \cos\theta}{\cos\theta - \cos\theta_0} d\theta = \frac{\pi \sin n\theta_0}{\sin\theta_0} \right) \rightarrow$$



## < 4.7 Classical Thin Airfoil Theory >

### ❖ The Symmetric Airfoil

$$\gamma(\theta) = 2\alpha V_\infty \left( \frac{1 + \cos\theta}{\sin\theta} \right)$$

\* Check Kutta condition

$$\lim_{\theta \rightarrow \pi} \gamma(\theta) = 2\alpha V_\infty \frac{0}{0} \quad \leftarrow \text{Indeterminant form}$$

By L'Hospital's rule

$$\begin{aligned} \rightarrow \gamma(\theta) &= \lim_{\theta \rightarrow \pi} 2\alpha V_\infty \left( \frac{-\sin\theta}{\cos\theta} \right)_{\theta=\pi} \\ &= 2\alpha V_\infty (0) = 0 \end{aligned}$$

## < 4.7 Classical Thin Airfoil Theory >

### ❖ The Symmetric Airfoil

\* Since we get  $\gamma(\theta)$ , now calculate  $\Gamma$ ,  $L$

$$\rightarrow \Gamma = \int_0^c \gamma(\xi) d\xi = \int_0^\pi \frac{c}{2} \gamma(\theta) \sin\theta d\theta \rightarrow \Gamma = \alpha c V_\infty \int_0^\pi (1 + \cos\theta) d\theta = \pi \alpha c V_\infty$$

\* **Lift** :  $L' = \rho_\infty V_\infty \Gamma = \pi \alpha c \rho_\infty V_\infty^2$

\* **Lift coefficient** :  $C_l = \frac{L'}{q_\infty c} = \frac{\pi \alpha c \rho_\infty V_\infty^2}{\frac{1}{2} \rho_\infty V_\infty^2 c} = 2\pi \alpha$

\* **Lift slope** :  $\frac{dC_l}{d\alpha} = 2\pi$

→ Lift coefficient is linearly proportional to angle of attack.

## < 4.7 Classical Thin Airfoil Theory >

### ❖ The Symmetric Airfoil

\* The moment about the leading edge

$$\begin{aligned}\rightarrow M_{LE}' &= - \int_0^c \xi (dL) \\ &= - \rho_\infty V_\infty \int_0^c \xi \gamma(\xi) d\xi \\ &= - \rho_\infty V_\infty \int_0^\pi \frac{c}{2} (1 - \cos\theta) 2\alpha V_\infty \left( \frac{1 + \cos\theta}{\sin\theta} \right) \frac{c}{2} \sin\theta d\theta \\ &= - \rho_\infty V_\infty^2 (2\alpha) \left( \frac{c}{2} \right)^2 \int_0^\pi (1 - \cos^2\theta) d\theta \\ &= - \frac{1}{2} \rho_\infty V_\infty^2 \alpha c^2 \int_0^\pi \sin^2\theta d\theta \\ &= - q_\infty c^2 \frac{\pi\alpha}{2}\end{aligned}$$



## < 4.7 Classical Thin Airfoil Theory >

### ❖ The Symmetric Airfoil

\* The moment coefficient

$$\rightarrow C_{m,LE} = -\frac{M_{LE}'}{q_{\infty} S c} = -\frac{M_{LE}'}{q_{\infty} c^2} = -\frac{\pi \alpha}{2}$$

$$C_l = 2\pi\alpha$$

$$C_{m,LE} = -\frac{C_l}{4}$$

$$C_{m,c/4} = C_{m,LE} + \frac{C_l}{4} = 0$$

\* Aerodynamic center is located at c/4 for incompressible, inviscid and symmetric airfoil (true in real world)

\* Center of pressure : the point at which the moment is zero

Aerodynamic center : the point at which the moment is independent of aoa

# Incompressible Flow over Airfoils

## < 4.8 The Cambered Airfoil >

\* From thin airfoil theory,

$$\rightarrow \frac{1}{2\pi} \int_0^c \frac{\gamma(\xi) d\xi}{(x-\xi)} = V_\infty \left( \alpha - \frac{dz}{dx} \right) \quad \dots (a)$$

\* For cambered airfoil,  $\frac{dz}{dx} \neq 0$

Transform  $\xi$  into  $\theta$   $\rightarrow \frac{1}{2\pi} \int_0^\pi \frac{\gamma(\theta) \sin \theta d\theta}{\cos \theta - \cos \theta_0} = V_\infty \left( \alpha - \frac{dz}{dx} \right) \quad \dots (b)$

\* The solution becomes

$$\gamma(\theta) = 2 V_\infty \left( A_0 \frac{1 + \cos \theta}{\sin \theta} + \sum_{n=1}^{\infty} A_n \sin n\theta \right) \quad \dots (c)$$

Leading term for symmetric airfoil

Fourier series term due to camber

## < 4.8 The Cambered Airfoil >

\* Substitute (c) into (b)

$$\rightarrow \frac{1}{\pi} \int_0^{\pi} \frac{A_0(1 + \cos\theta)d\theta}{\cos\theta - \cos\theta_0} + \frac{1}{\pi} \int_0^{\pi} \frac{A_n \sin n\theta \sin\theta d\theta}{\cos\theta - \cos\theta_0} = \alpha - \frac{dz}{dx}$$

By using the integral standard form

$$\int_0^{\pi} \frac{\sin n\theta \sin\theta d\theta}{\cos\theta - \cos\theta_0} = -\pi \cos n\theta_0$$

$$\int_0^{\pi} \frac{\cos n\theta d\theta}{\cos\theta - \cos\theta_0} = \frac{\pi \sin n\theta_0}{\sin\theta_0}$$

$$\rightarrow A_0 - \sum_{n=1}^{\infty} A_n \cos n\theta_0 = \alpha - \frac{dz}{dx}$$

$$\rightarrow \frac{dz}{dx} = (\alpha - A_0) + \sum_{n=1}^{\infty} A_n \cos n\theta_0$$

# Incompressible Flow over Airfoils

## < 4.8 The Cambered Airfoil >

For Fourier cosine series,

$$f(\theta) = B_0 + \sum_{n=1}^{\infty} B_n \cos n\theta, \quad 0 \leq \theta \leq \pi$$

$$B_0 = \frac{1}{\pi} \int_0^{\pi} f(\theta) d\theta$$

$$B_n = \frac{2}{\pi} \int_0^{\pi} f(\theta) \cos n\theta d\theta$$

$$\begin{aligned} \rightarrow A_0 &= \alpha - \frac{1}{\pi} \int_0^{\pi} \frac{dz}{dx} d\theta_0 \\ A_n &= \frac{2}{\pi} \int_0^{\pi} \frac{dz}{dx} \cos n\theta_0 d\theta_0 \end{aligned}$$

[Note] given  $\frac{dz}{dx}, \alpha \rightarrow$  Determine  $\gamma(\theta)$  to make the camber line a streamline with  $A_0, A_n$  + Kutta condition,  $\gamma(\pi)=0$

# Incompressible Flow over Airfoils

## < 4.8 The Cambered Airfoil >

\* The total circulation due to the entire vortex sheet

$$\begin{aligned}\Gamma &= \int_0^c \gamma(\xi) d\xi = \frac{c}{2} \int_0^\pi \gamma(\theta) \sin\theta d\theta \quad \leftarrow \gamma(\theta) = 2V_\infty \left( A_0 \frac{1 + \cos\theta}{\sin\theta} + \sum_{n=1}^{\infty} A_n \sin n\theta \right) \\ &= cV_\infty \left[ A_0 \int_0^\pi (1 + \cos\theta) d\theta + \sum_{n=1}^{\infty} A_n \int_0^\pi \sin n\theta \sin\theta d\theta \right]\end{aligned}$$

By using  $\int_0^\pi (1 + \cos\theta) d\theta = \pi$  ,  $\int_0^\pi \sin n\theta \sin\theta d\theta = \begin{cases} \frac{\pi}{2} & \text{for } n = 1 \\ 0 & \text{for } n \neq 1 \end{cases}$

$$\rightarrow \Gamma = cV_\infty \left( \pi A_0 + \frac{\pi}{2} A_1 \right)$$



## < 4.8 The Cambered Airfoil >

\* Lift coefficient for a cambered thin airfoil

$$\rightarrow L' = \rho_{\infty} V_{\infty} \Gamma = \rho_{\infty} V_{\infty}^2 c \left( \pi A_0 + \frac{\pi}{2} A_1 \right)$$

$$\rightarrow C_l = \frac{L'}{\frac{1}{2} \rho_{\infty} V_{\infty}^2 c} = \pi (2A_0 + A_1)$$

$$\begin{cases} A_0 = \alpha - \frac{1}{\pi} \int_0^{\pi} \frac{dz}{dx} d\theta_0 \\ A_1 = \frac{2}{\pi} \int_0^{\pi} \frac{dz}{dx} \cos \theta_0 d\theta_0 \end{cases}$$

$$\rightarrow C_l = 2\pi \left( \alpha + \frac{1}{\pi} \int_0^{\pi} \frac{dz}{dx} (\cos \theta_0 - 1) d\theta_0 \right)$$

$$\rightarrow \text{Lift slope, } a_0 = \frac{dC_l}{d\alpha} = 2\pi$$

# Incompressible Flow over Airfoils

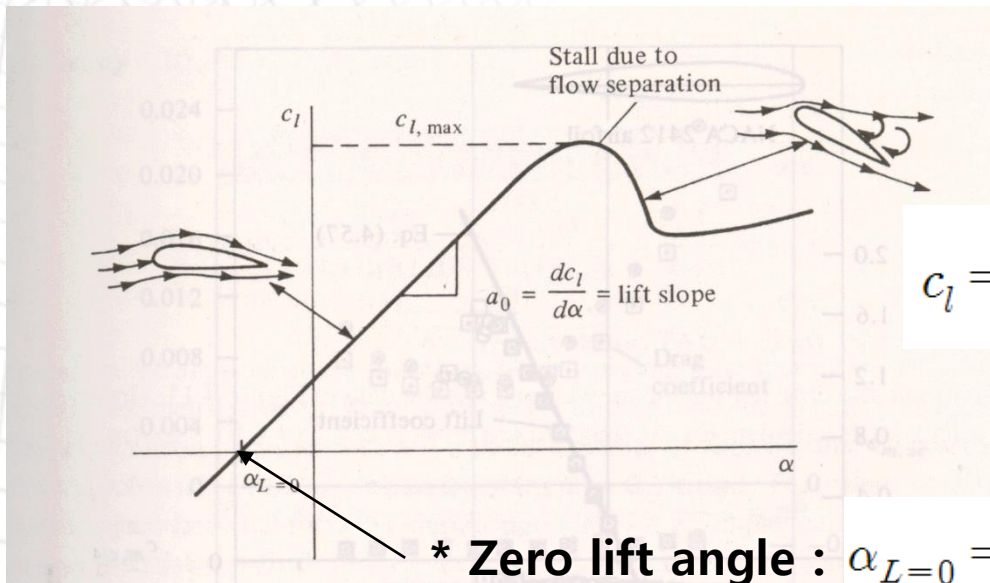
## < 4.8 The Cambered Airfoil >

\* Lift coefficient for a cambered thin airfoil

$$C_l = 2\pi\left(\alpha + \frac{1}{\pi} \int_0^\pi \frac{dz}{dx} (\cos\theta_0 - 1) d\theta_0\right)$$

[Note]

\* Lift slope is  $2\pi$  for any shape airfoil



$$c_l = \frac{dc_l}{d\alpha} (\alpha - \alpha_{L=0}) = 2\pi (\alpha - \alpha_{L=0})$$

\* Zero lift angle :  $\alpha_{L=0} = -\frac{1}{\pi} \int_0^\pi \frac{dz}{dx} (\cos\theta_0 - 1) d\theta_0$

## < 4.8 The Cambered Airfoil >

\* The total moment about the leading edge

$$\begin{aligned}
 M_{LE} &= - \int_0^c \xi(dL) = - \int_0^c \rho_\infty V_\infty \gamma(\xi) d\xi \\
 &= - \rho_\infty V_\infty \int_0^\pi 2 V_\infty \left( A_0 \frac{1 + \cos\theta}{\sin\theta} + \sum_{n=1}^{\infty} A_n \sin n\theta \right) \frac{c}{2} (1 - \cos\theta) \left( \frac{c}{2} \sin\theta \right) d\theta \\
 &= - \frac{1}{2} \rho_\infty V_\infty^2 c^2 \left[ \int_0^\pi A_0 (1 - \cos^2\theta) d\theta + \int_0^\pi A_n \sin n\theta \sin\theta d\theta - \int_0^\pi A_n \sin n\theta \sin\theta \cos\theta d\theta \right] \\
 &= - \frac{1}{2} \rho_\infty V_\infty^2 c^2 \left[ \frac{\pi}{2} \left( A_0 + A_1 - \frac{A_2}{2} \right) \right]
 \end{aligned}$$

\* Moment coefficient

$$\begin{aligned}
 C_{m,LE} &= \frac{M_{LE}}{q_\infty c^2} = - \frac{\pi}{2} \left( A_0 + A_1 - \frac{A_2}{2} \right) \quad \leftarrow C_l = \pi(2A_0 + A_1) \\
 &= - \left[ \frac{C_l}{4} + \frac{\pi}{4} (A_1 - A_2) \right]
 \end{aligned}$$

$$\rightarrow C_{m, \frac{c}{4}} = \frac{\pi}{4} (A_2 - A_1)$$

→  $A_1$  &  $A_2$  both are independent of  $\alpha$   
 → The quarter-chord is the aerodynamic center for a cambered airfoil

## < 4.8 The Cambered Airfoil >

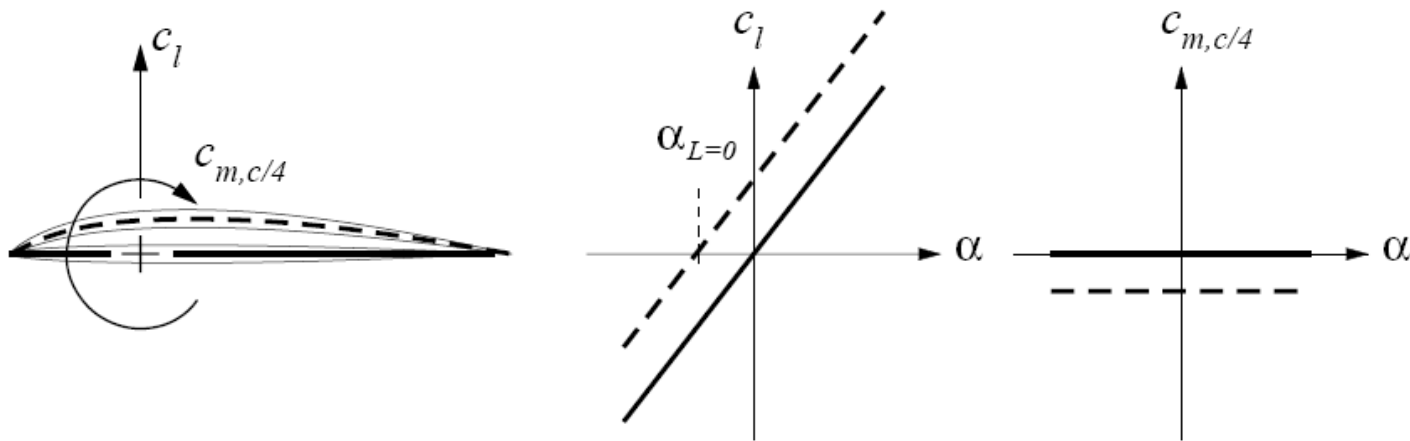
\* The center of pressure

$$\begin{aligned}x_{cp} &= -\frac{M_{LE}}{L} = -\frac{C_{m,LE}c}{C_l} \\ &= \frac{\left[\frac{C_l}{4} + \frac{\pi}{4}(A_1 - A_2)\right]c}{C_l} \\ &= \frac{c}{4}\left[1 + \frac{\pi}{C_l}(A_1 - A_2)\right] \rightarrow f(\alpha) \rightarrow \text{Not a convenient point}\end{aligned}$$

# Incompressible Flow over Airfoils

## < 4.8 The Cambered Airfoil >

### ❖ The influence of camber on the thin airfoil



---\* The cambered airfoil

$$C_l = 2\pi\left(\alpha + \frac{1}{\pi} \int_0^\pi \frac{dz}{dx} (\cos\theta_0 - 1) d\theta_0\right)$$

$$C_{m,\frac{c}{4}} = \frac{\pi}{4}(A_2 - A_1)$$

——\* The symmetric airfoil

$$C_l = 2\pi\alpha$$

$$C_{m,c/4} = C_{m,LE} + \frac{C_l}{4} = 0$$